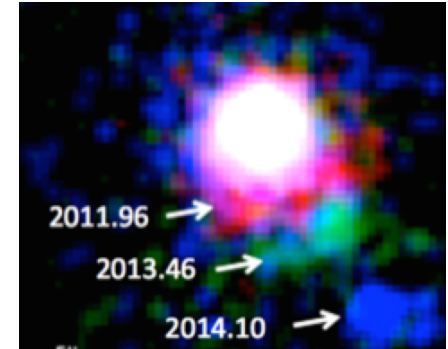


Extended X-ray object ejected from the PSR B1259-63/LS 2883 binary



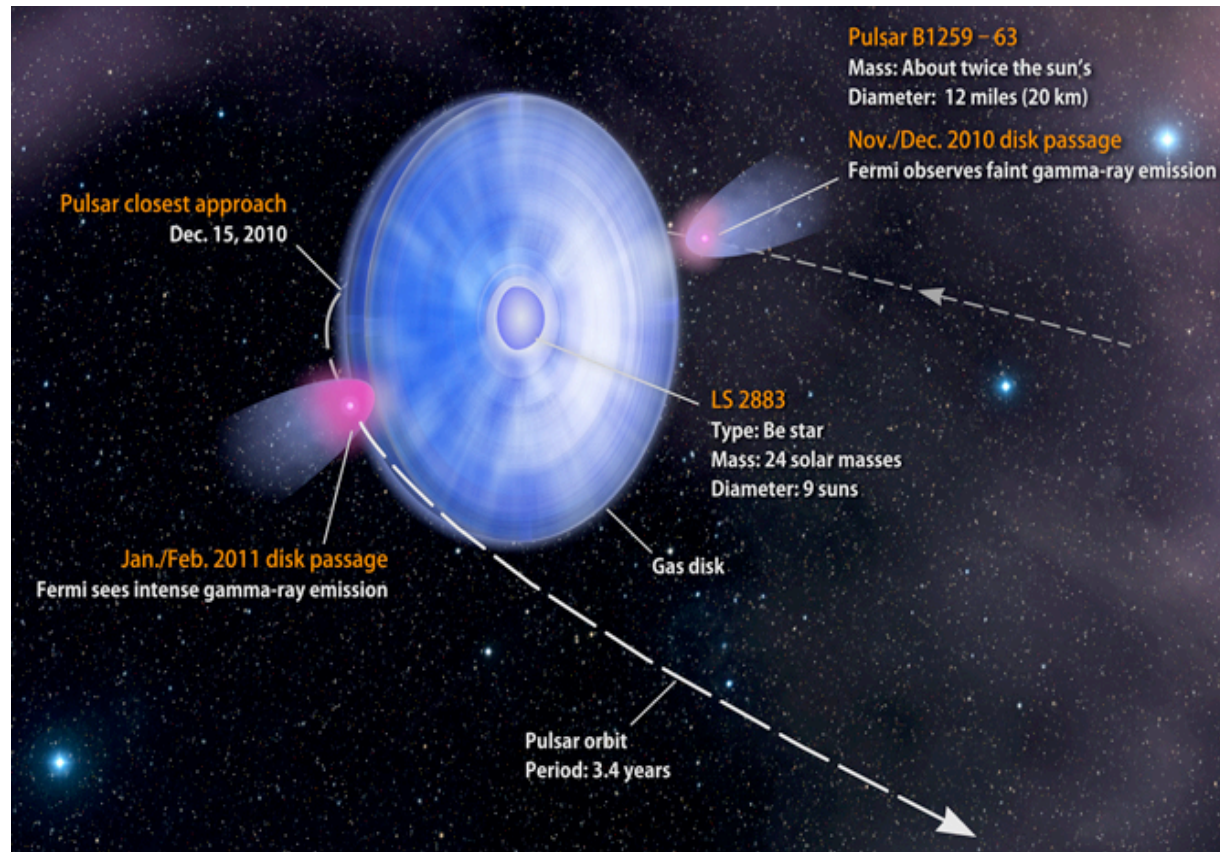
Oleg Kargaltsev (George Washington University)

George Pavlov (Pennsylvania State University)

Jeremy Hare (George Washington University)

Blagoy Rangelov (George Washington University)

High-mass binary LS 2883 with PSR B1259-63



(Credit: NASA's Goddard Space Flight Center/Francis Reddy)

Fast-spinning, massive ($M \sim 30 M_{\odot}$, $L = 6 \times 10^4 L_{\odot}$) star with a strong wind.

The **wind** is dense and slow in the **decretion disk**, tenuous and fast outside the disk.

Pulsar B1259-63:

Spin period = 48 ms

$\dot{E} = 8 \times 10^{35}$ erg/s

Spin-down age = 330 kyr

Should emit pulsar wind

X-ray flux varies with orbital period.
Gamma-ray flashes near periastron, apparently when the pulsar interacts with the decretion disk during 2nd passage.

Orbit:

3.4 yr orbital period

7 AU (3 milliarcsec) max. separation

0.87 eccentricity

Imaging observations with Chandra ACIS

4 observations, May 2009 – Feb 2014

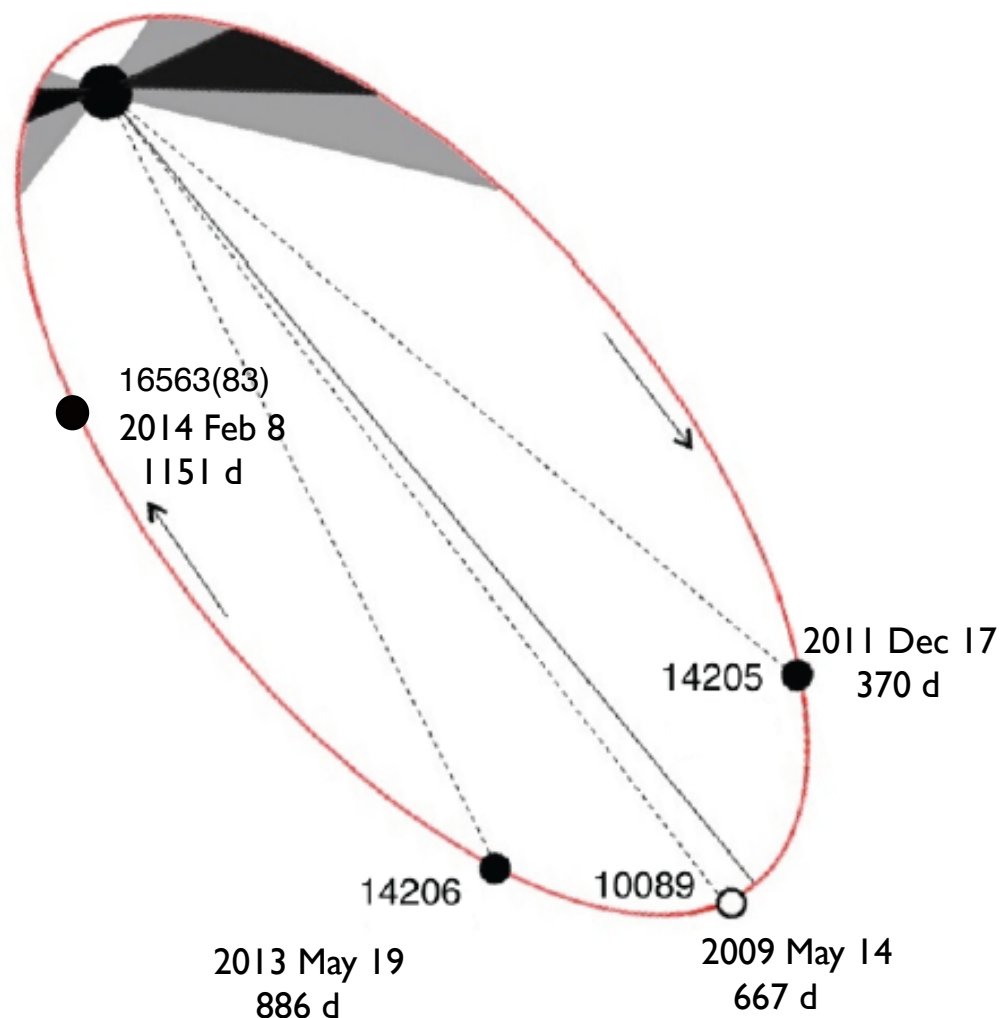
ObsID	MJD	θ^a deg	Δt^b days	Exp. ^c ks	Cts ^d
10089	54965	182	667	25.6	1825 61
14205	55912	169	370	56.3	6551 343
14206	56431	192	886	56.3	4162 144
<i>new</i> ⁱ	56696	221	1151	57.6	6257 58

^aTrue anomaly counted from periastron.

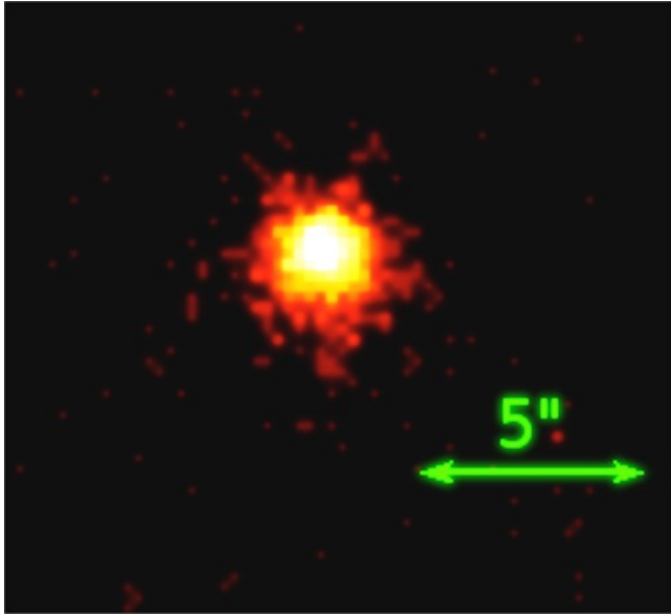
^bDays since latest preceded periastron.

^cExposure corrected for deadtime.

^dTotal (gross) counts.

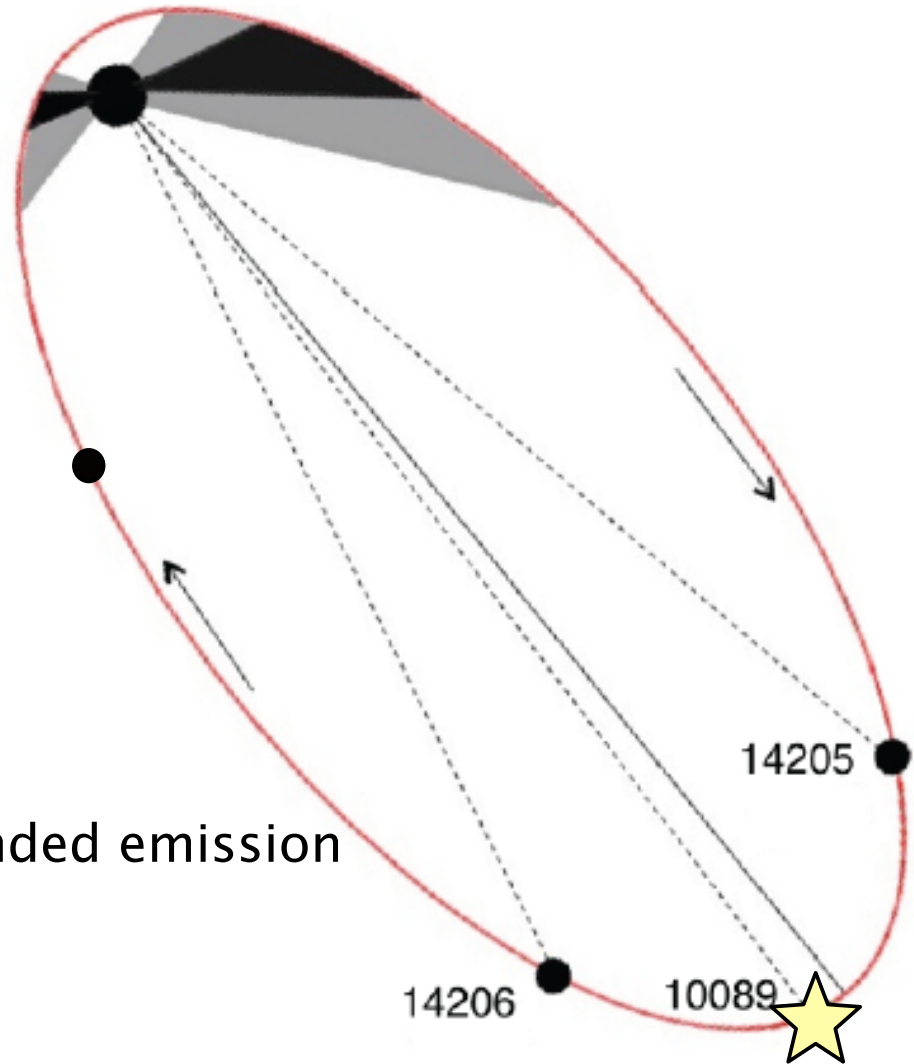


1st Observation (2009 May 14)

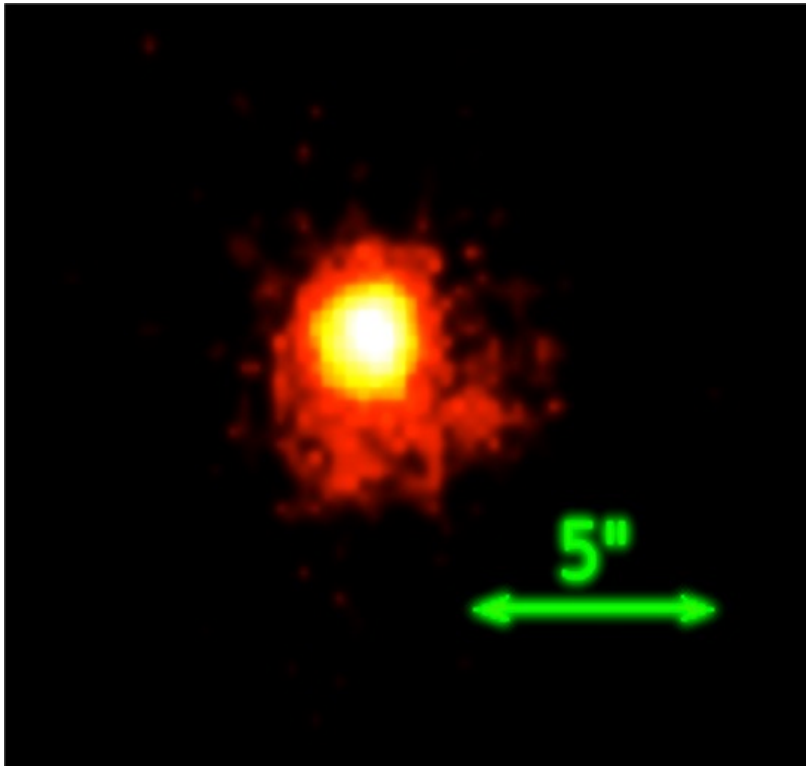


Short 25.6 ks ACIS-I exposure
near apastron, $\theta = 182$ deg

$\sim 4\sigma$ detection of asymmetric extended emission
(Pavlov et al 2011)

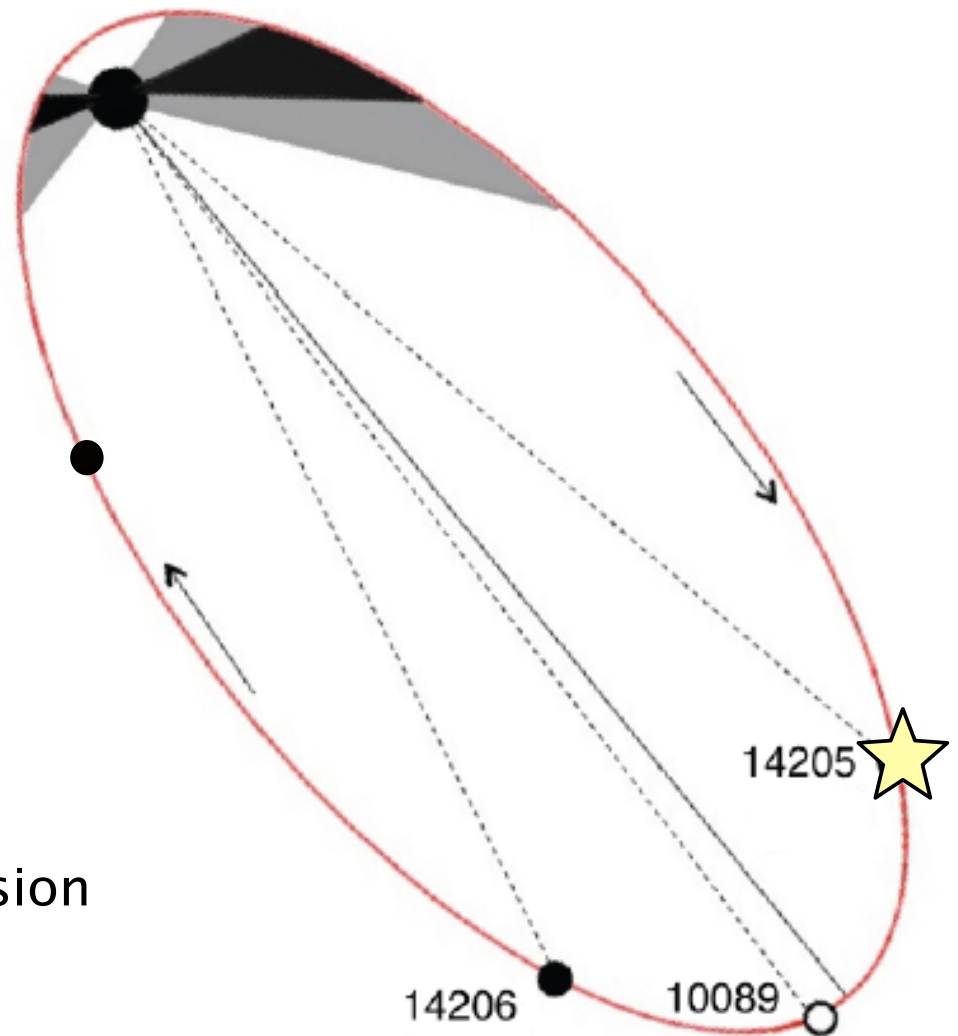


2nd Observation (2011 Dec 17)

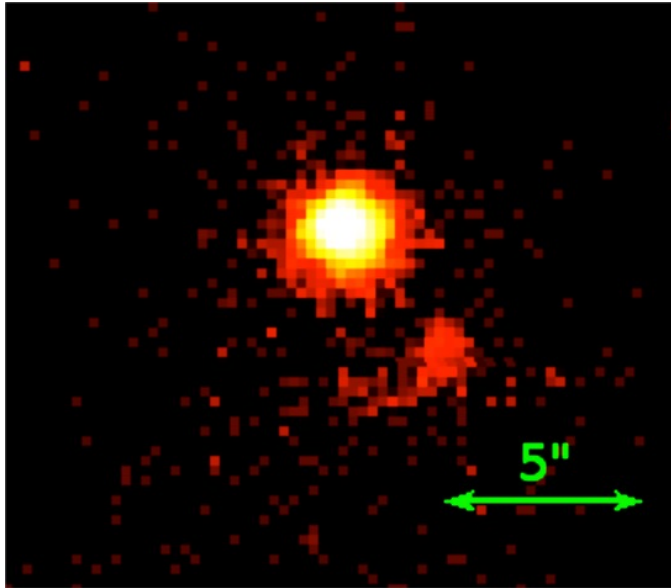


56.3 ks ACIS-I exposure
before apastron, $\theta = 169$ deg

Clear asymmetric extended emission

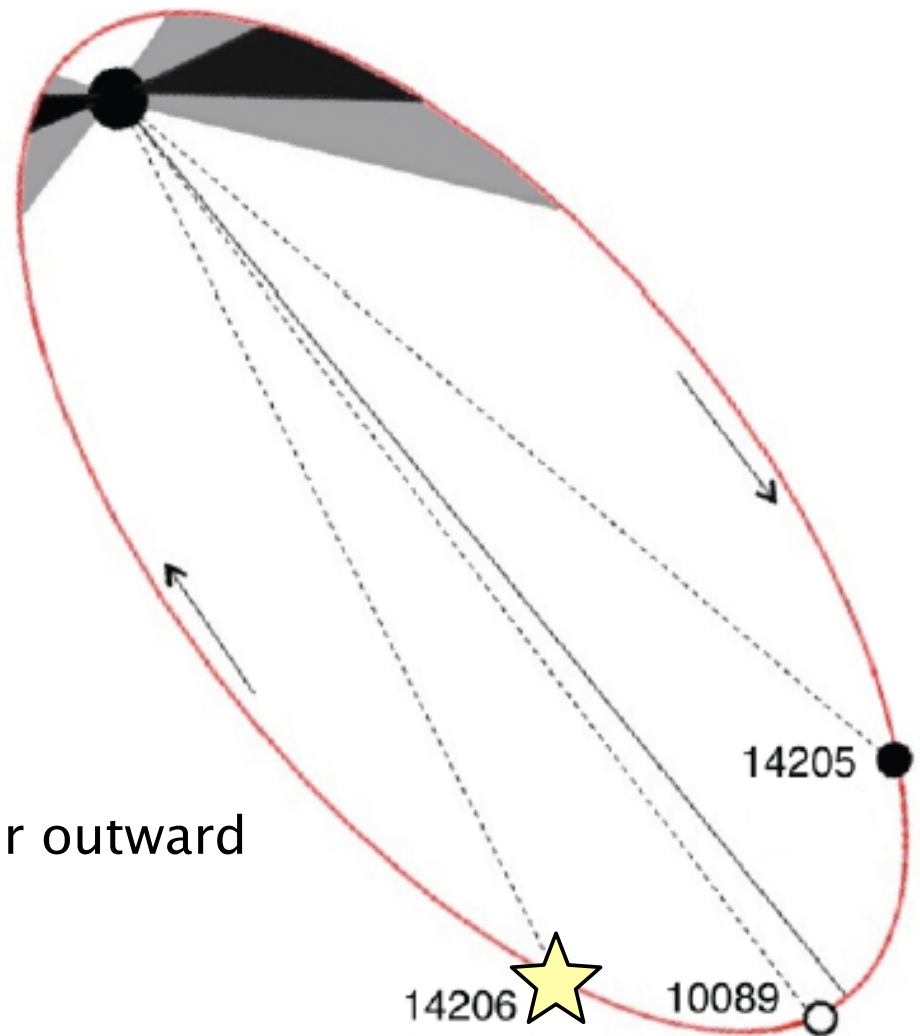


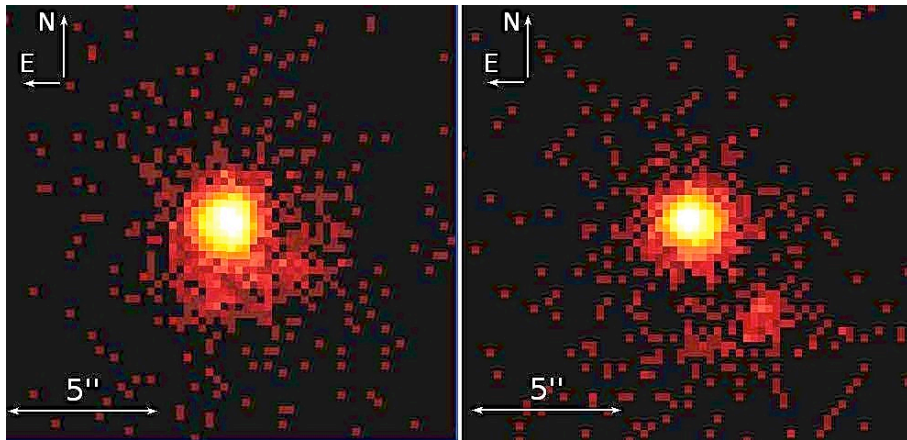
3rd Observation (2013 May 19)



56.3 ks ACIS-I exposure
after apastron, $\theta = 192$ deg

Extended emission moved further outward





2nd and 3rd observations compared

$1''.8 \pm 0''.5$ shift

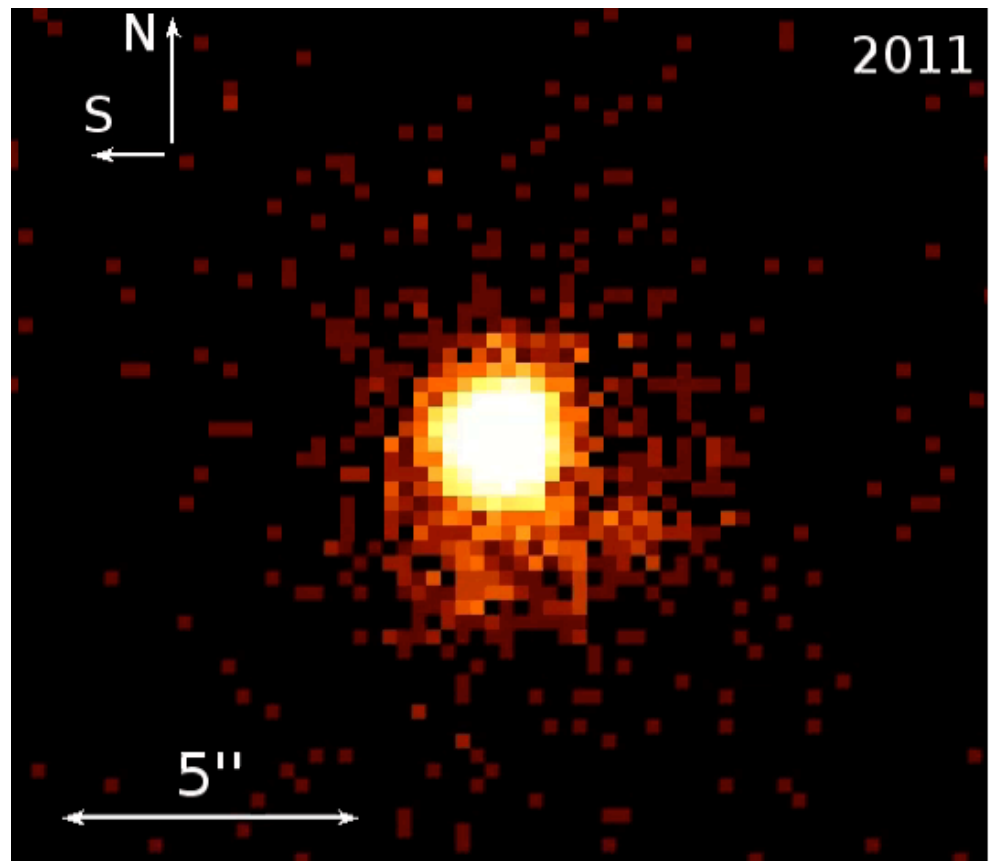
corresponds to the
apparent proper motion

$$\mu = 1.27 \pm 0.35 \text{ arcsec yr}^{-1}$$

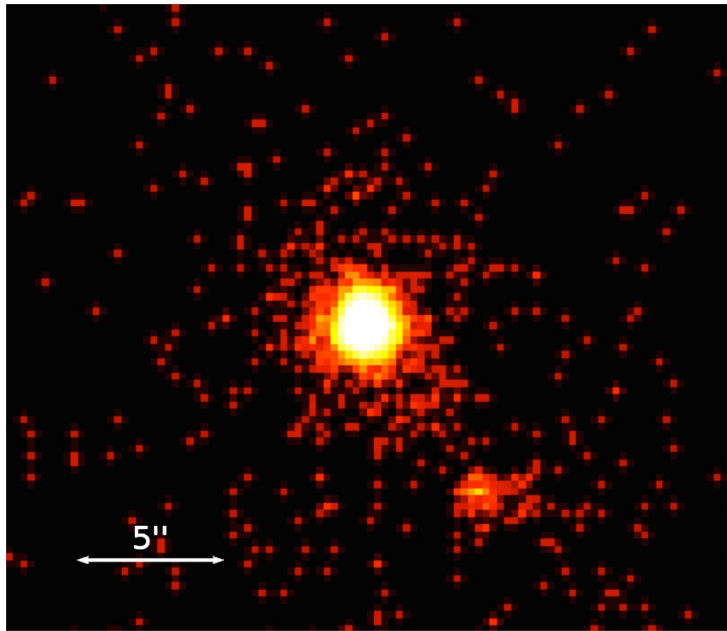
$$\mathbf{V} = (0.046 \pm 0.013)\mathbf{c}$$

at $d = 2.3 \text{ kpc}$

(Kargaltsev et al. 2014)

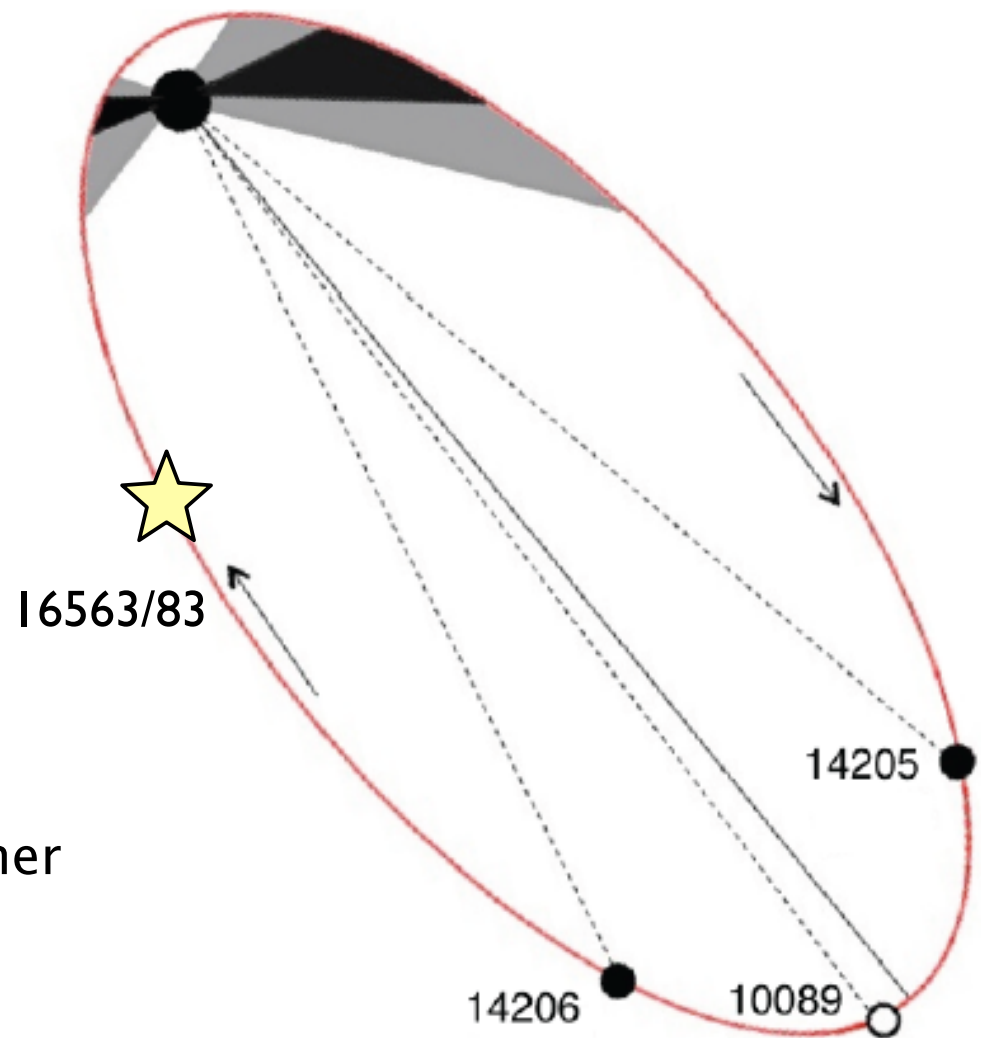


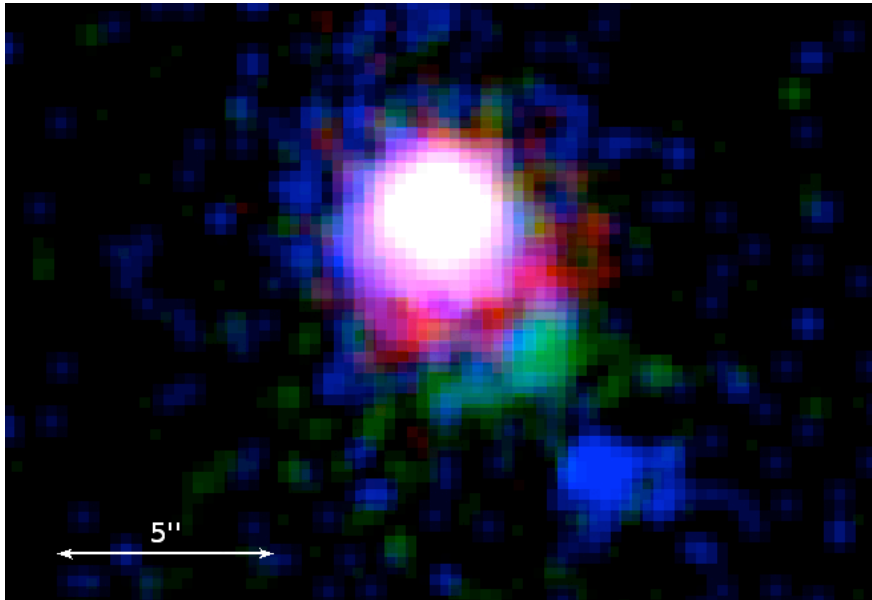
4th Observation (2014 February 8 - 9)



57 ks ACIS-I exposure
approaching periastron

Extended emission moved farther
from the binary, apparently
faster than expected from the
previous 2 observations





Between 3rd and 4th observations
the extended structure moved by
 $2.5'' \pm 0.5''$.

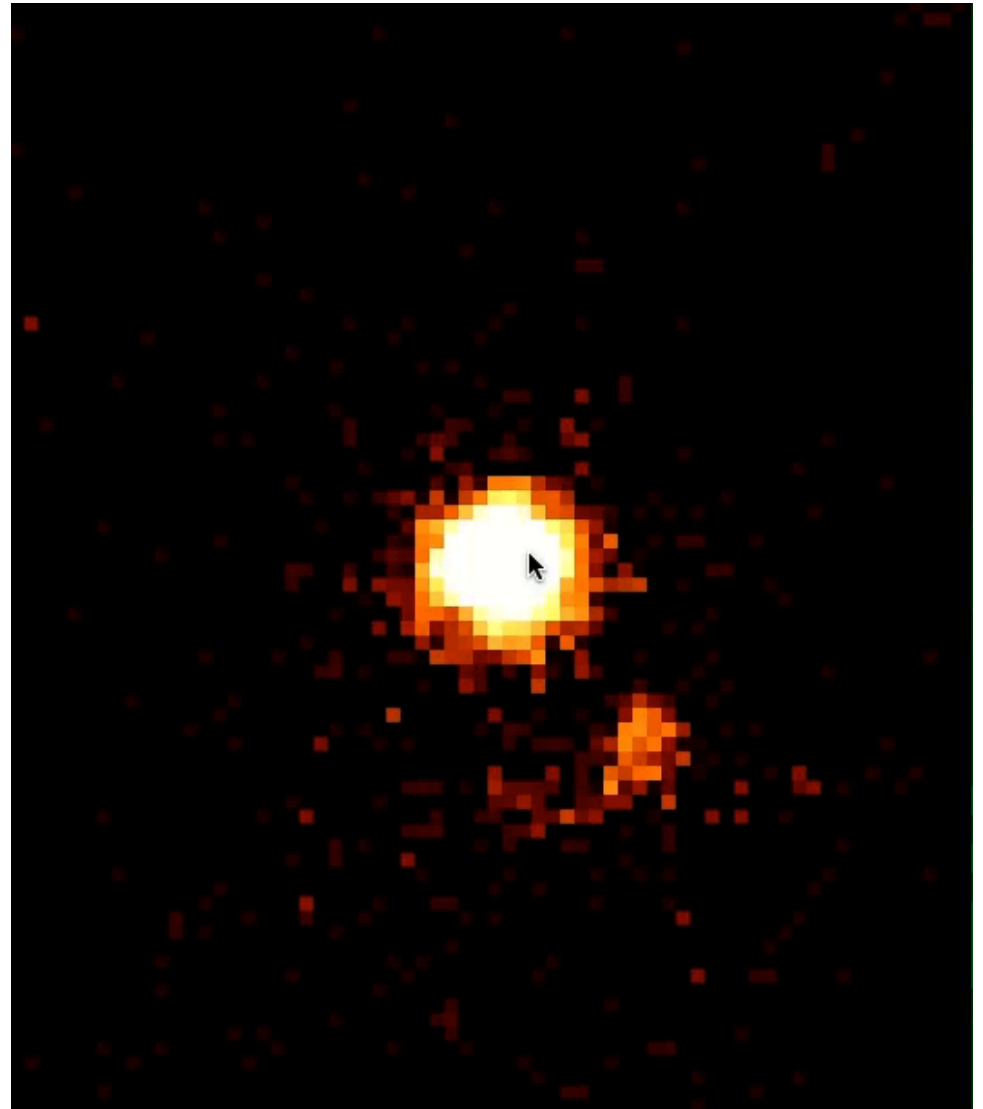
This corresponds to the apparent
proper motion

$$V = (0.13 \pm 0.03)c$$

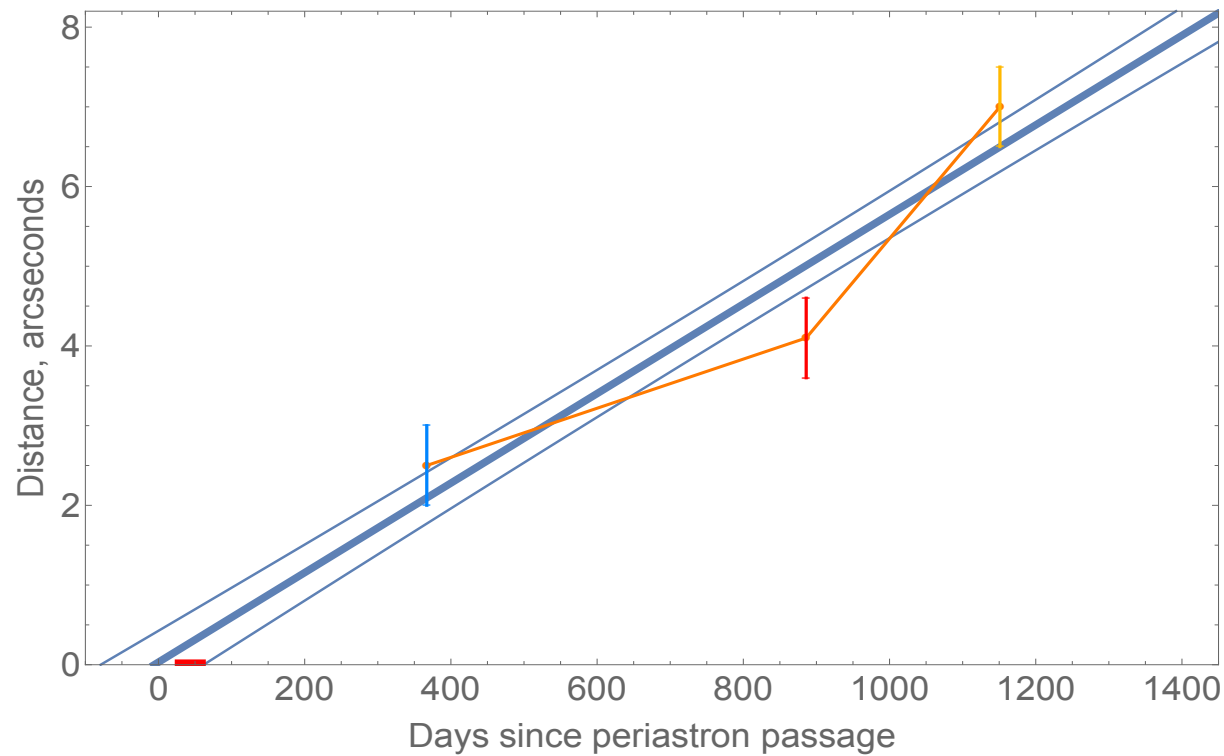
at $d = 2.3 \text{ kpc}$

Apparent **acceleration** (?)
 $90 \pm 40 \text{ cm s}^{-2}$

2nd, 3rd, and 4th observations
together:



Distance of the extended source from the binary versus time



Linear fit: $V = (0.07 \pm 0.01)c$

If there is no (or little) acceleration, the cloud was ejected from the binary around periastron of 2010 Dec 14

Luminosities and spectra of extended emission

In 3 last observations 0.5 – 8 keV fluxes are

$$\mathbf{F = 8.5+/-0.5, \quad 3.6+/-0.4, \quad 1.9+/-0.4 \quad \times 10^{-14} \text{ erg cm}^{-2} \text{ s}^{-1} ,}$$

corresponding luminosities $\mathbf{L \sim (0.2 - 1) \times 10^{31} \text{ erg/s}}$ at $d = 2.3 \text{ kpc}$,
~0.7% - 3% of the binary's luminosity.

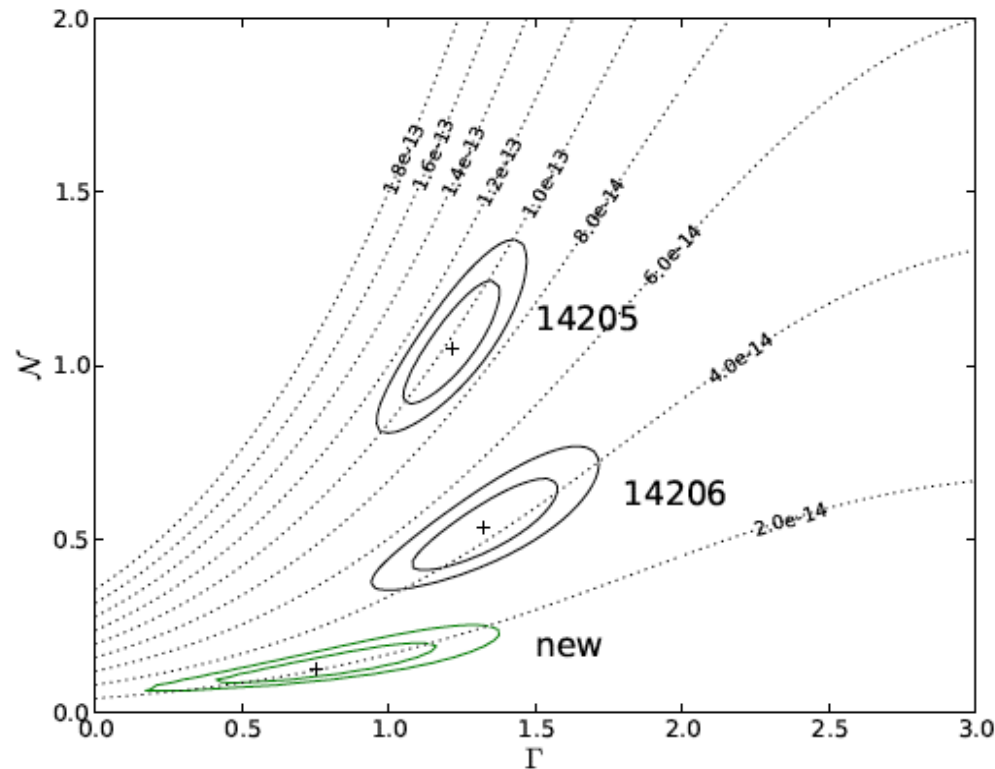
The spectra can be fitted with thermal bremsstrahlung,
 $kT > 6 \text{ keV}$, $n \sim 100 \text{ cm}^{-3}$, $\mathbf{m_{ej} \sim 10^{28} - 10^{29} \text{ g}}$ -- much larger than
the mass supplied by the massive star wind during one orbital period,
 $P_{\text{orb}} \dot{M} \sim 10^{26} (\dot{M}/10^{-8} M_{\text{sol}}/\text{yr}) \text{ g}$, or a reasonable mass of disk,
 $m_{\text{disk}} \sim 10^{24} - 10^{26} \text{ g}$.
Kinetic energy $\sim 10^{46} - 10^{47} \text{ erg}$, improbably large.

**The scenario with hot hadronic plasma cloud radiating
via bremsstrahlung does not look plausible.**

The spectra are also consistent with power laws,
photon indices $\Gamma = 1.2 \pm 0.1$, 1.3 ± 0.2 , and 0.8 ± 0.4 (no softening!)

Synchrotron radiation?

Confidence contours in Photon Index – Normalization plane



Synchrotron interpretation:

magnetic field $\mathbf{B} \sim 80 \mathbf{k}_m^{2/7} \mu\mathbf{G}$, where $k_m = \epsilon_{\text{mag}}/\epsilon_{\text{kin}}$;

electron Lorentz factor $\gamma \sim 10^7 - 10^8$,

total magnetic and electron energies $\mathbf{W}_m \sim 5 \times 10^{40} \mathbf{k}_m^{4/7}$ and $\mathbf{W}_e \sim 5 \times 10^{40} \mathbf{k}_m^{-3/7} \text{ erg}$ in volume $\mathbf{V} \sim 10^{50} \text{ cm}^3$.

$\mathbf{W}_m + \mathbf{W}_e \ll P_{\text{orb}} \dot{E} = 9 \times 10^{43} \text{ erg}$ for a broad range of k_m -- the energy could be supplied by the pulsar.

But, if the ejected object were an **e-/e+ cloud**, it would be difficult to explain the fast motion because of the **drag force**, $\mathbf{f} \sim \rho_{\text{amb}} \mathbf{v}^2 \mathbf{A}$.

Deceleration time

$$\mathbf{t}_{\text{dec}} \sim (\mathbf{W}_m + \mathbf{W}_e) \mathbf{v} \mathbf{f}^{-1} \mathbf{c}^{-2} \sim 10 \mathbf{n}_{\text{amb}}^{-1} (\mathbf{k}_m^{4/7} + \mathbf{k}_m^{-3/7}) \mathbf{s},$$

where n_{amb} is the ambient proton number density.

To reduce the deceleration, the e-/e+ cloud must be loaded with a heavy (electron-ion) plasma, but even in this case the ejected mass should be a substantial fraction of the disk mass, if the ejected clump is moving in a stellar wind blown bubble.

Another hypothesis

Variable termination shock in the circumbinary medium,

similar to PWNs around isolated pulsars (Kargaltsev et al 2014)

But, it requires unrealistically high ambient pressure, $p_{\text{amb}} \sim 10^{-10} \text{ dyn cm}^{-2}$, to explain the shock size; looks artificial now.

Current explanation: Instead of the companion's wind bubble,

ejected clump is moving in the unshocked pulsar wind

More plausible at larger values of $\eta = \dot{E}/(\dot{M} v_w c) =$
 $= 4.4 (\dot{M}/10^{-9} M_{\odot}/\text{yr})^{-1} (v_w/1000 \text{ km/s})^{-1}$

when the companion's wind is confined by the pulsar wind into a narrow cone, while the unshocked pulsar wind fills the rest of the binary volume.

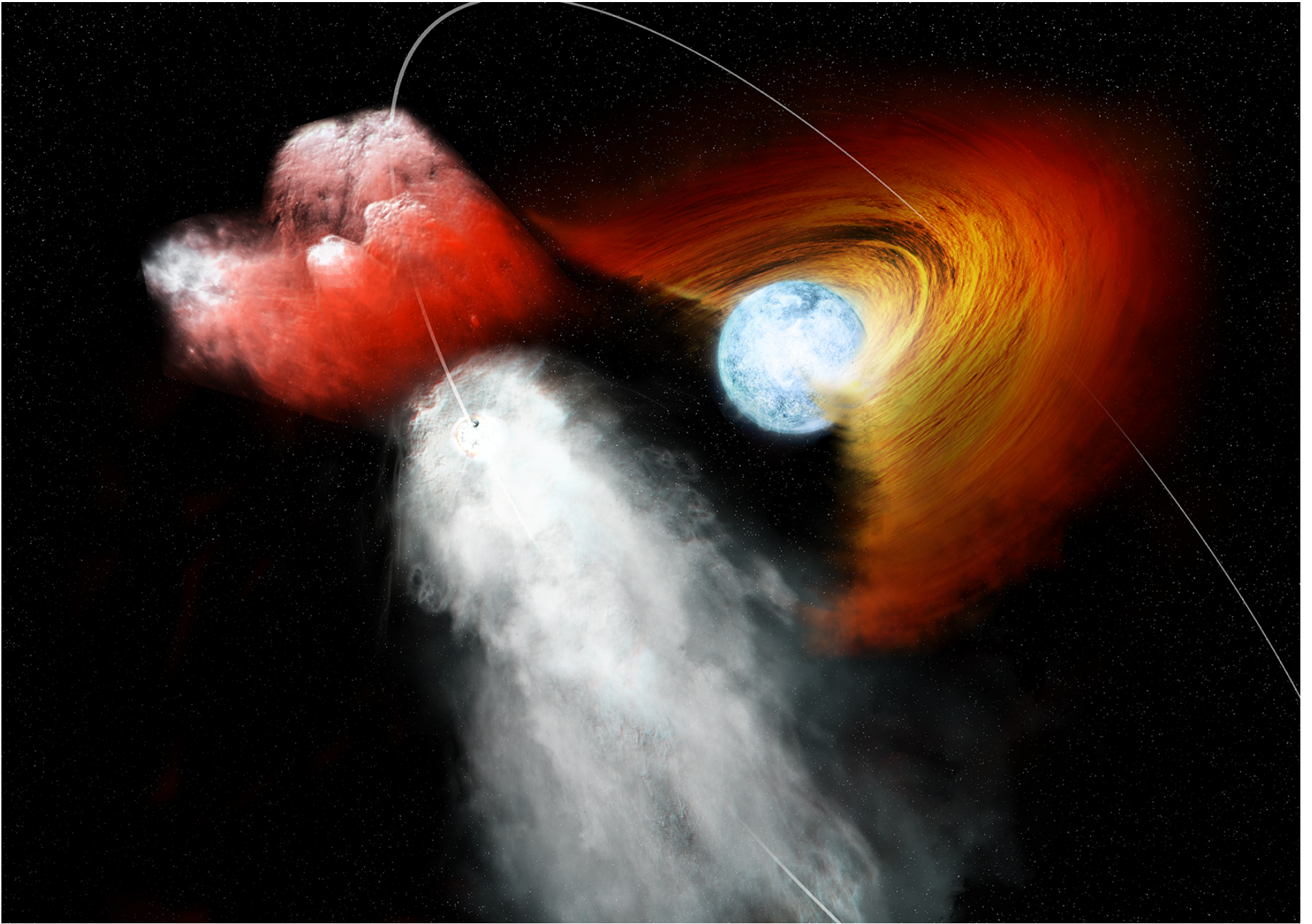
The X-ray emission is due to synchrotron radiation of the pulsar wind shocked by the collision with the clump.

X-ray luminosity $L_{\text{X,cl}} = \xi_X \dot{E} (r_{\text{cl}}/2r)^2$, $\xi_X \sim 1.5 \times 10^{-3}$

The interaction with unshocked pulsar wind with ejected clump can also **accelerate** the clump: $\dot{v} \sim p_{pw} A m_{cl}^{-1}$.

$m_{cl} \sim 10^{21} \text{ g}$ for the apparent (low-significance) estimated acceleration.

The clump could be ejected due to interaction of the pulsar with the decretion disk. When the pulsar enters the dense disk, a shock is created, with a radius exceeding the disk's vertical size →
→ Disruption of the disk in the first passage, further fragmentation in the second passage, γ -ray flares from shocked pulsar wind, entrainment of clumps in the pulsar wind, then acceleration by the pulsar wind ram pressure to $\sim 0.1 \text{ c}$.



Artist interpretation: NASA/CXC press release

Summary

- We discovered a new phenomenon: Ejection of an X-ray emitting clump from a high-mass γ -ray binary with a velocity $v \sim 0.1c$ and a hint of acceleration.
- The clump's luminosity faded with time but the power-law-like spectrum ($\Gamma \sim 0.8 - 1.3$) did not show softening.
- The clump was likely ejected due to interaction of the pulsar (pulsar wind) with the equatorial decretion disk of the high-mass star.
- We suggest that the clump is moving in the unshocked pulsar wind, whose pressure accelerated the clump to the very high speed. This scenario requires large η .
- The most likely emission mechanism is synchrotron radiation of relativistic electrons ($E_e \sim 10 - 100 \text{ TeV}$, $B \sim 10^2 \mu\text{G}$) of pulsar wind shocked in the collision with clump.
- We expect a new clump has been ejected during the recent periastron passage (May 2014), new Chandra observations are planned.